

Fiber-amplifier pumped high average power few-cycle pulse non-collinear OPCPA

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Abstract: We report on the performance of a 60 kHz repetition rate sub-10 fs, optical parametric chirped pulse amplifier system with 2 W average power and 3 GW peak power. This is to our knowledge the highest average power sub-10 fs kHz-amplifier system reported to date. The amplifier is conceived for applications at free electron laser facilities and is designed such to be scalable in energy and repetition rate.

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References and links

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1. Introduction

Ultra-short few-cycle-pulse laser amplifier systems have revolutionized many fields of research in physics. Ultra-short pulse generation based on optical parametric amplification [1] has found wide application in spectroscopy, where tunable sources are needed [2, 3]. It was shown that multi-millijoule few-cycle pulses can be attained at high repetition rate [4] or with high peak power [5, 6, 7]. A multi-TW peak power amplifier system based on the OPCPA technique has already been used for experiments in plasma physics [8]. A substantial overview on existing OPA and OPCPA systems is given in reference [9].

Parametric amplification combined with the chirped pulse amplification technique is to date the only way possible for amplification of sub-10 fs pulses at highest average power. A further increase in average power in OPCPA systems requires high average power pump lasers. Recently, fiber laser systems have reached the millijoule pulse energy range with femtosecond pulse durations [10] and average power as high as 830 W with femtosecond pulses [11, 12]. This generation of amplifier systems is well suited as OPCPA pump laser [13, 14]. Certainly, the tendency in OPCPA amplifier development will in future shift towards higher flux which is required for precision pump-probe experiments and coincidence measurements. Also the latest generation of storage ring-based synchrotron radiation sources and linear accelerator-based free-electron laser (FEL) sources [15] require this kind of laser amplifiers especially for pump probe experiments.

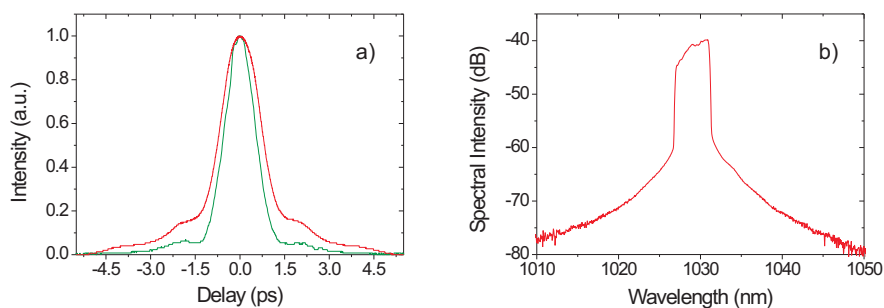


Fig. 2. (a) Autocorrelation measurement of the fundamental pump amplifier pulses (red) and autocorrelation measurement of the second harmonic (green); (b) Spectrum of the amplified 1030 nm pulses.

spectrum and the autocorrelation measurement ($\tau_{AC}=1.6$ ps, thus $\tau_{FWHM} \sim 1.1$ ps assuming a gaussian pulse shape) are shown in Fig. 2(a) and (b). This radiation is frequency doubled in a 1 mm type I BBO crystal with an efficiency of 60% and used to pump the OPA stage (330 μ J, 20 W, 515 nm). The pulse duration of the second harmonic is also measured Fig. 2(a) and amounts to $\tau_{FWHM}=840$ fs.

The other 40% of the Ti:Sapphire pulse energy is used to seed a non-collinear optical parametric amplifier stage. A spatial light modulator (SLM) in 4-f geometry is used for adaptive control of the spectral phase [20]. This prism based SLM setup is used in combination with a prism stretcher to maximize the transmission through the shaping and stretcher setup ($\eta_{stretcher+SLM} \sim 50\%$). The prism separation is chosen such to obtain a net negative dispersion of second, third and fourth order, to be used in combination with a glass compressor. The pulses from the Ti:Sapphire oscillator are stretched to 520 fs FWHM to match the temporal pump pulse gain window. The negatively chirped pulses sustain already compression in the 4 mm BBO crystal (OPA stage) down to 300 fs FWHM. The remaining dispersion is compensated in a 10 mm fused silica alignment-free compressor ($\eta_{compressor} = 95\%$) and by the phase modulator (SLM) up to the sixth order. The CPA setup is configured such to have a compensated second order dispersion. The calculated residual spectral phase coefficients $\frac{\partial^i \phi}{\partial \omega^i}$ of order $i > 2^{nd}$ to 6^{th} at a central wavelength of 760 nm are $-1.58E3$ fs³, $-1.18E4$ fs⁴, $1.10E4$ fs⁵ and $-5.26E4$ fs⁶. These dispersion orders are compensated with the spatial light modulator.

The OPA stage is currently working in the tangential phase matching scheme at a non-collinear angle of 2.22° and the phase matching angle at 24.15° . Angular detuning allows to amplify a larger bandwidth in a 4 mm long type I BBO crystal (see Fig. 3a, dotted line) compared to the ideal phase-matched case. The pump and seed beam diameter used are 0.9 mm and 0.75 mm (FWHM), respectively. The pump intensity is 25 GW/cm² and the resulting gain is on the order of $>10^4$.

3. Amplifier performance

The 0.5 nJ seed beam is amplified to 35 μ J of single pulse energy. The average power from seed pulses which are not amplified in the OPA stage is already taken into account. The amplified signal power versus the pump pulse power is shown in Fig. 3(b). The maximum average power extraction is reached at the saturation point of this stage. The level of amplified optical parametric fluorescence is 30 mW as determined by blocking the seed beam and measuring the

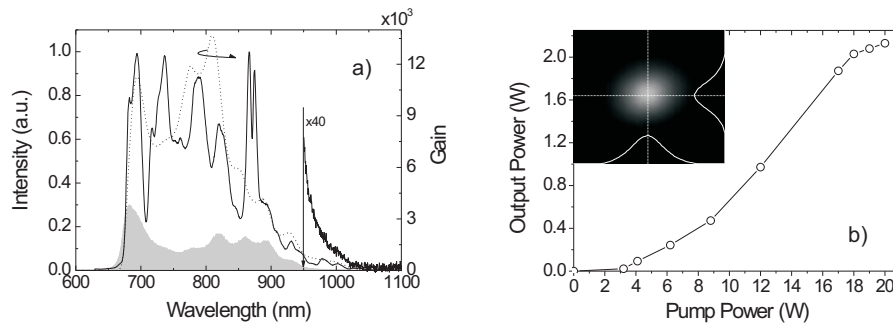


Fig. 3. (a) Spectrum of the 35 uJ, 60 kHz amplified signal pulses (solid line) compared to simulation of gain (dotted line) and oscillator spectrum (shaded grey and solid black line 40x magnification from 950-1100nm); (b) Measured energetics of the optical parametric amplifier stage and of the (inset) amplified beam profile.

OPF content after a ~ 3 mm aperture at 1 meter distance from the OPA stage. The effective optical parametric fluorescence (OPF) present in an area which equals the FWHM of the seed beam is smaller by more than one order of magnitude, which is a consequence of the different divergence of the signal beam compared to OPF. Additionally, an OPF quenching effect is expected if the signal beam is present and amplified. An estimate is given by numerical simulation [21] which gives a quenching factor of 3. This results in a value of OPF smaller than 1 mW in presence of the seed beam and taking into account the different divergence of the signal and OPF beam. In this case, the temporal contrast ratio after compression is better than 5 orders of magnitude.

At this saturation level of the OPA stage, also conversion in the second harmonic has been observed. 10% of the signal output was converted, leaving a total amplified signal average output power of 2 W taking into account the compressor efficiency or $\sim 33.3 \mu\text{J}$ compressed pulse energy. The spectrum of the amplified signal is shown in Fig. 3(a) and its corresponding beam profile in the inset of Fig. 3(b). The pulse duration measurement in this regime of amplification turned out to be challenging due to the large amount of second harmonic light present. A second order autocorrelation measurement is performed with a commercial autocorrelator (Femto-meter, Femtolasers GmbH). The correlation trace in Fig. 4(a) shows a deconvoluted compressed pulse duration of 7.8 fs at FWHM (assuming a gaussian pulse shape) attaining a pulse peak power of 3 GW by estimating the energy content at FWHM.

The gain was reduced by a factor of ~ 2 to reduce generation of the second harmonic of the signal beam (1.2 W output power) to be able to perform more accurate pulse duration measurements. The compressed pulse duration for these parameters is measured via autocorrelation and is 7 fs. A SPIDER measurement was performed to confirm this result and gives a pulse duration measurement of 6.9 fs. There is still some room of improvement, the Fourier limited pulse duration amounts to ~ 6 fs. In this regime of amplification, the amount of amplified optical parametric fluorescence is substantially lower and the average power stability of the OPCPA is 1.3% rms. The pulse stability was measured in unit of power, with an acquisition rate of 4 kHz. The compression results for the $33.3 \mu\text{J}$ pulses are shown in Fig. 4(a) while the compression results for the $20 \mu\text{J}$ pulses are shown in Fig. 4(b)-(d).

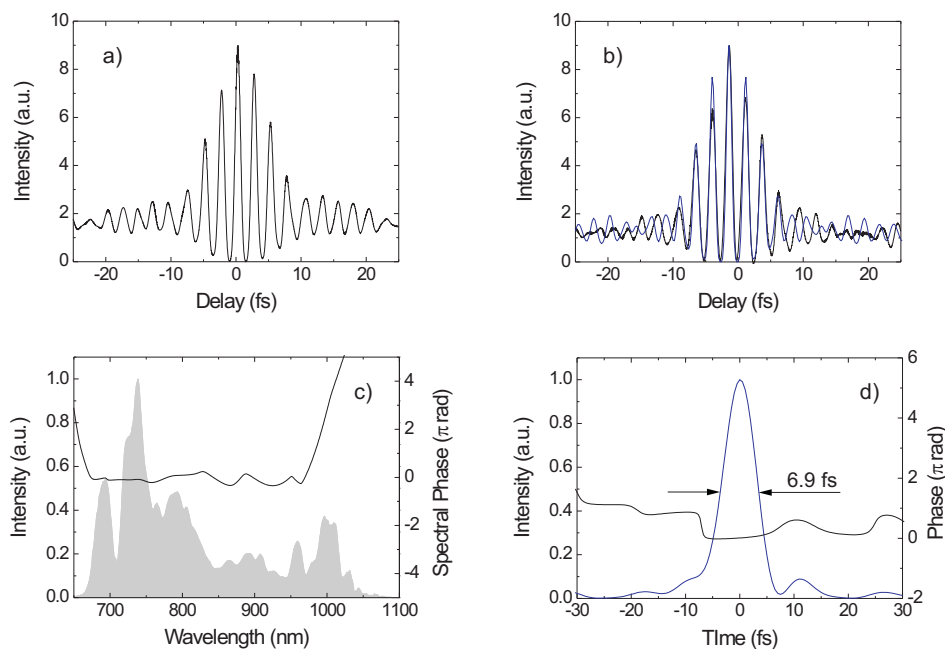


Fig. 4. (a) Interferometric autocorrelation trace of 33.3 μJ pulses b) interferometric autocorrelation trace of 20 μJ signal pulses (black) compared to the calculated autocorrelation trace from a SPIDER measurement (blue). (c and d) SPIDER measurement of the same 20 μJ signal pulses: (c) spectral phase and amplified spectrum; (d) temporal phase and reconstructed temporal pulse shape.

4. Conclusions

We have reported on the generation of 6.9 fs pulses with 20 μJ and 7.8 fs pulses with 33.3 μJ single pulse energy, at 60 kHz repetition rate from an ultrabroadband optical parametric chirped pulse amplifier system. The pulses are stable within 1.3% rms and the calculated maximal pulse peak power amounts to 2.03 and 3 GW, respectively. This optical parametric chirped pulse amplifier represents the first step towards the development of a new generation of ultrashort-pulse high repetition rate amplifiers to be deployed at linear accelerator-based free electron laser EUV and X-ray sources. The pump amplifier technology is scalable in energy and repetition rate [12, 22]. The final objective is to reach MHz repetition rate and several millijoule output pulse energy. Moreover, the output of the optical parametric amplifier has the potential to be shortened in pulse duration. Carrier envelope offset stabilization of the amplified pulses is as well possible and can be implemented in future upgrades of the amplifier.

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